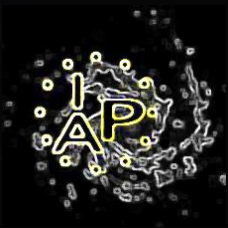


Looking for ultra-high energy astroparticles in a radio haystack



Simon Chiche – Institut d'Astrophysique de Paris

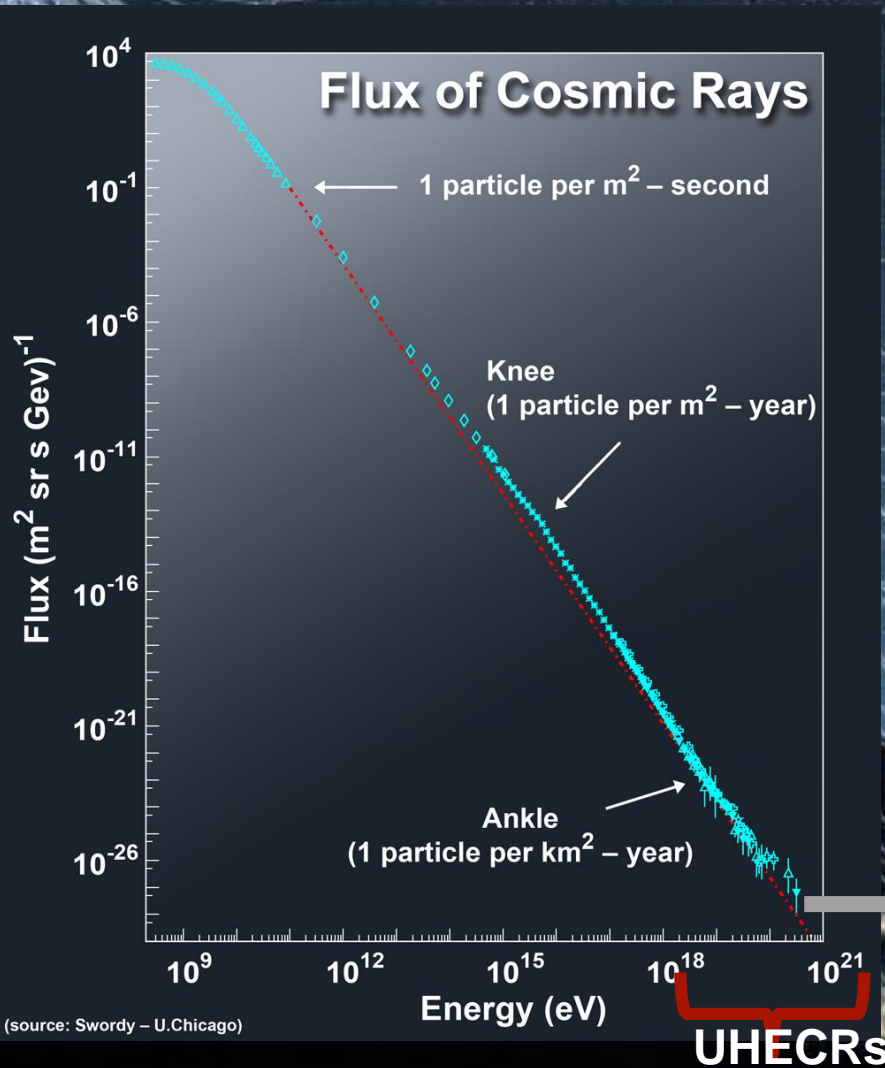
Supervisors: Kumiko Kotera (IAP), Olivier Martineau (LPNHE)



The mystery of ultra-high energy cosmic rays (UHECRs)

- Cosmic rays: high energy atomic nuclei (protons, iron nuclei, etc)
- Most energetic particles in the universe (ultra-high energy cosmic rays: $E > 10^{18} eV$)
- **Where do they come from?**

- At the lowest energy: Solar origin
- Intermediate energy: SNR (galactic origin)
- **Ultra-high energy: ?**



We don't know the exact nature of these particles

We don't know the sources

We don't know the acceleration mechanisms

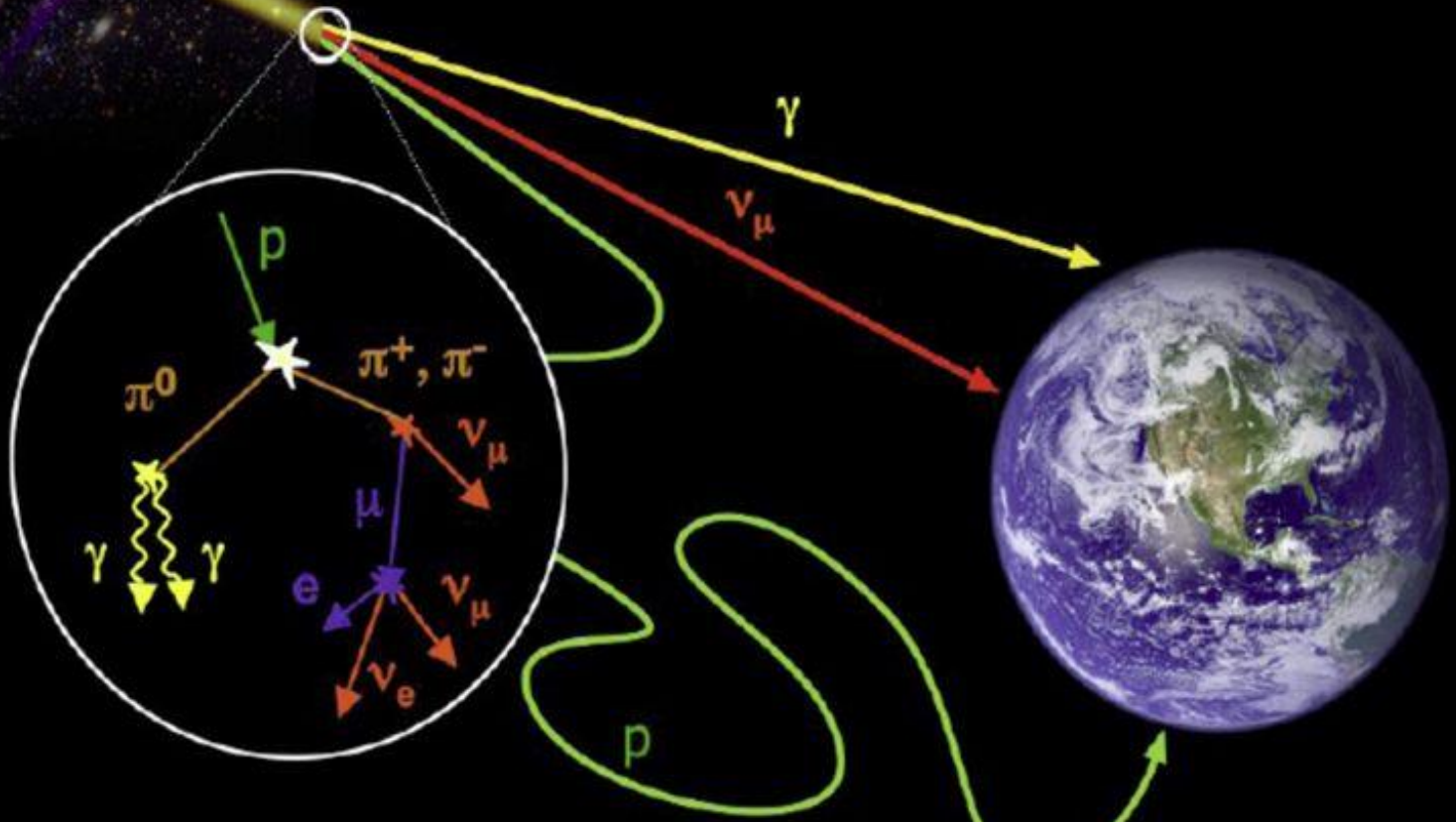
Very low flux:
 $1. km^{-2}. century^{-1}$

Ultra-high energy multi-messengers (UHE)!

- ✓ probe the most powerful sources in the Universe
- ✓ understand the origin of ultra-high energy cosmic rays

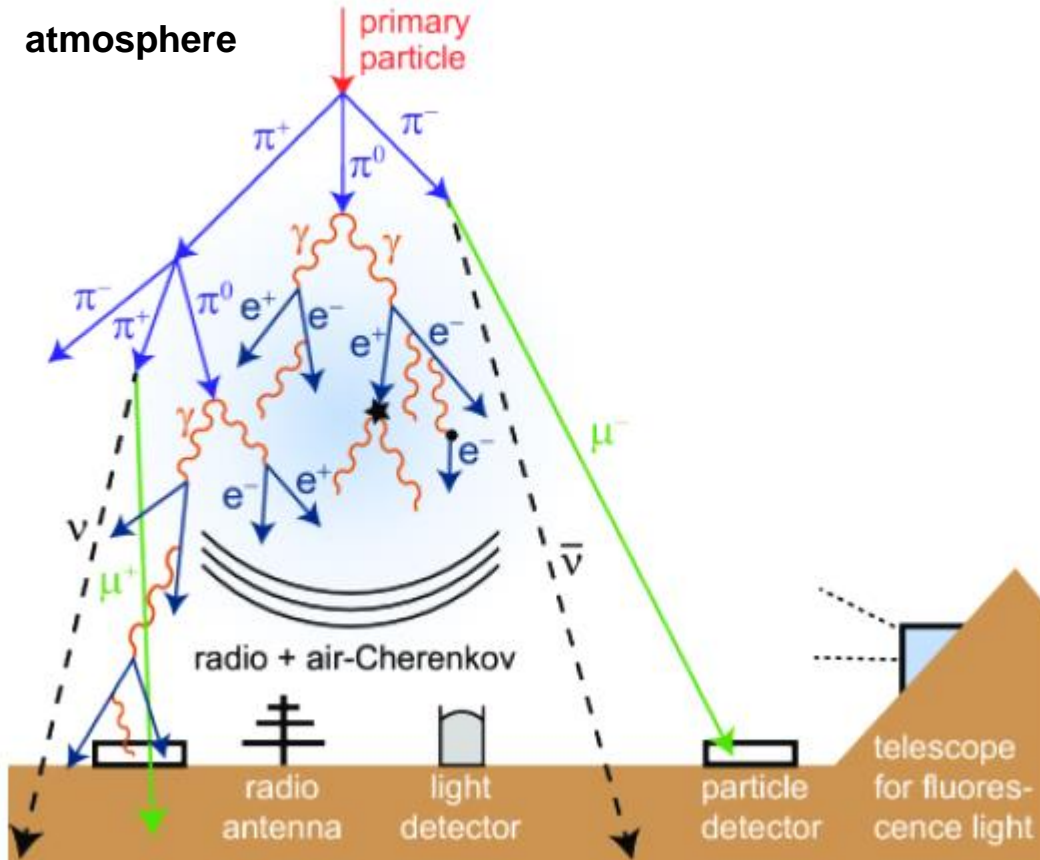
?

+ Gravitational waves



Extensive air showers (EAS)

Interaction of high energy astroparticles with the atmosphere: shower/cascade of secondary particles!



- Hadronic component: mainly π decaying into μ and ν
- Electromagnetic part: e^+ , e^- , γ

Main emissions:

- Cherenkov light
- Fluorescence light
- Radio emission

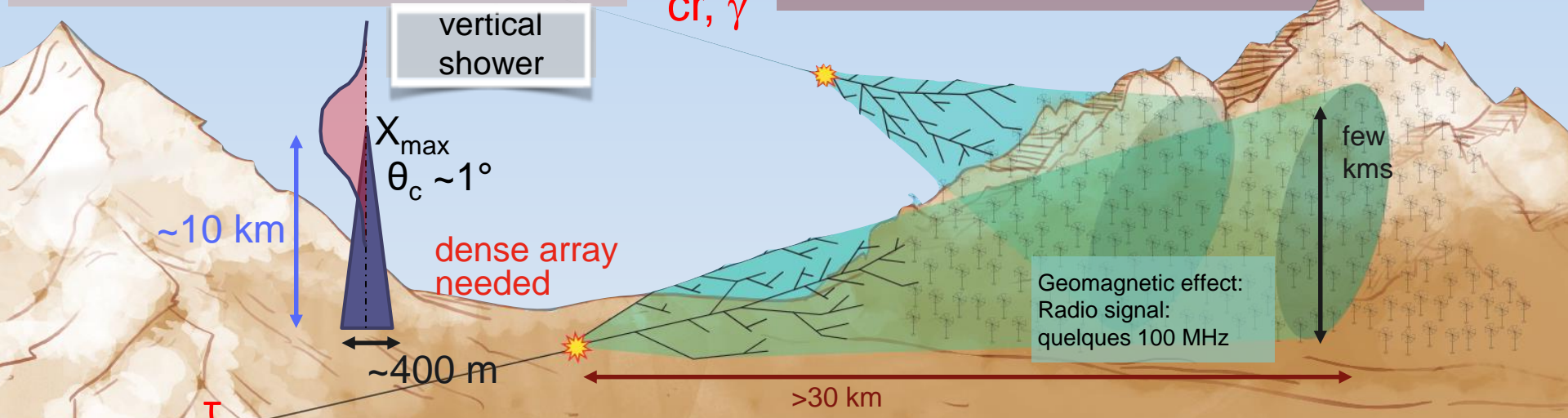
➔ We can detect the signal originating from the electromagnetic part with radio antennas!

GRAND and GRANDproto300

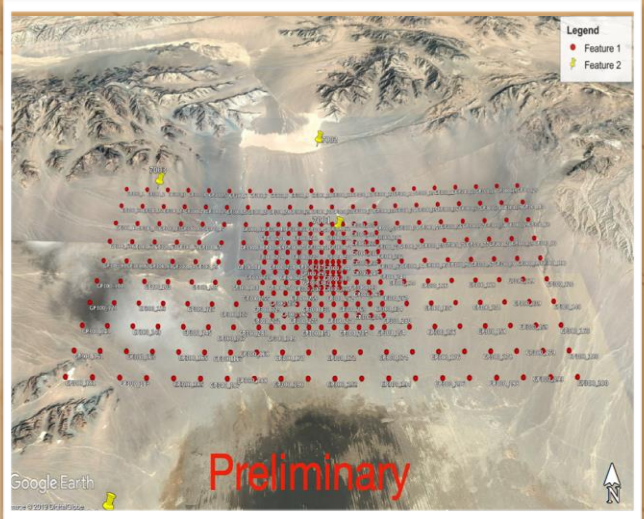
GRAND : Giant radio array of 200 000 radio antennas over 200 000 km^2

Detection of air showers induced by ultra-high energy astroparticles

Inclined showers with mountains as targets



GRANDProto300
first prototype in
2021!

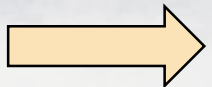


Prototype of 300 antennas, 200 km^2

Detection of astroparticles with $E_{range} = 10^{16.5} - 10^{18}\text{ eV}$

GRANDProto300: Challenges of radiodetection

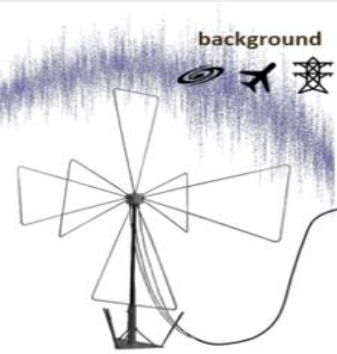
- Autonomous detection of astroparticles



Grail of radiodetection!

Current experiments: external triggers (Cerenkov tanks, scintillators)

GRAND: radio antennas only



Overwhelming noise from human emissions

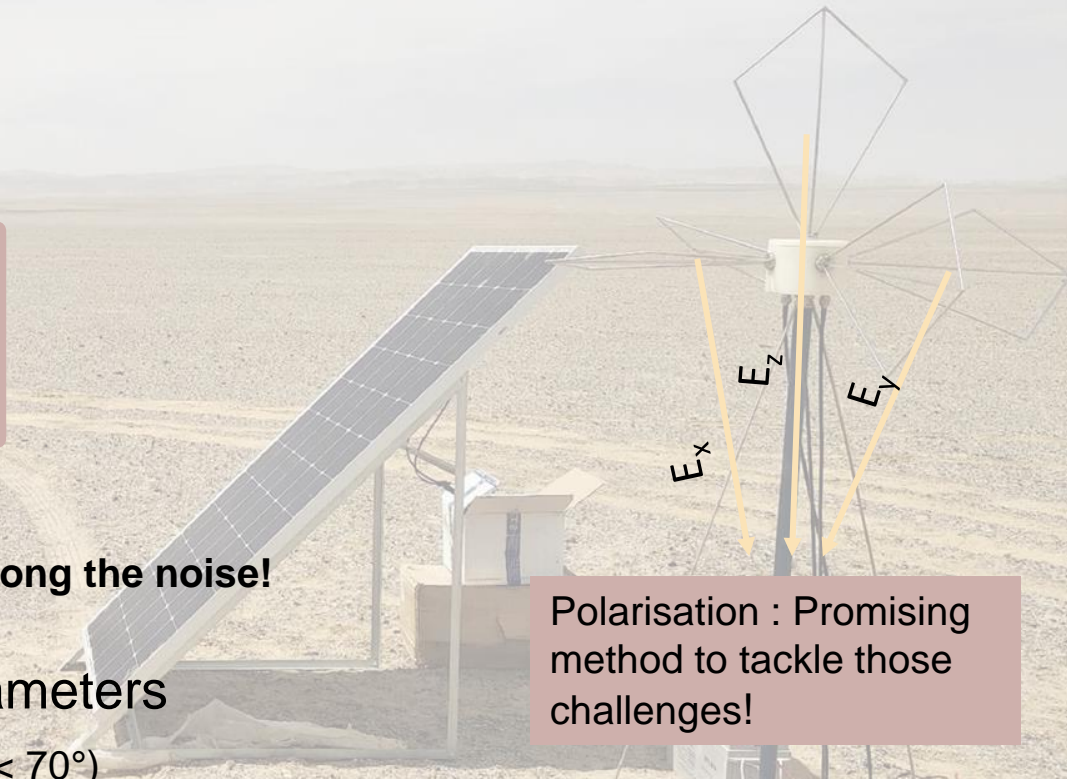
We have to identify the radio signal among the noise!

- Reconstruction of shower parameters

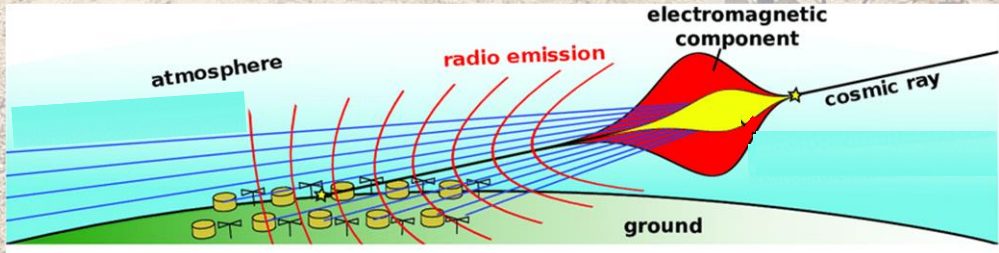
Current experiments: vertical showers ($\theta < 70^\circ$)

GRAND detection of inclined showers ($\theta > 70^\circ$)

Asymmetries, ground reflections effects



Polarisation : Promising method to tackle those challenges!

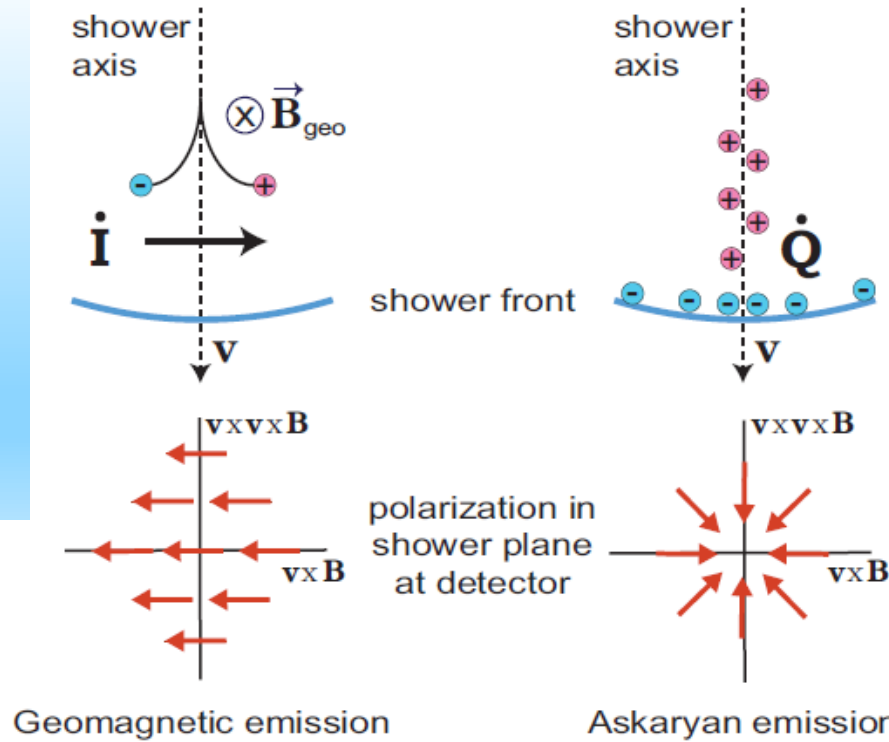


Polarisation of the radio signal

Polarisation: direction of the electric field

Geomagnetic emission

- Induced dipole with \vec{B}_{geo}
- Polarisation along $-\vec{v} \times \vec{B}$
- Main contribution to the radio signal



Charge excess emission

- Accumulation of negative charges close to the shower core
- Radial polarisation
- $\approx 10\%$ of the amplitude of the total emission for vertical air showers

Schröder (2017)

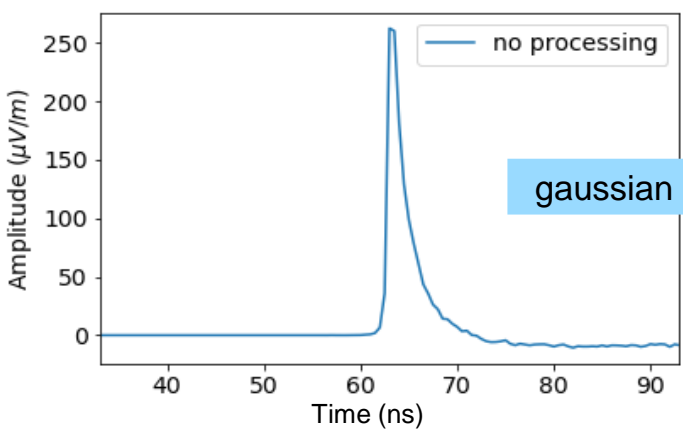
- Complex polarisation signature: allows to discriminate the signal from the noise
- Charge excess signature: gives insights about the core position

Traces processing

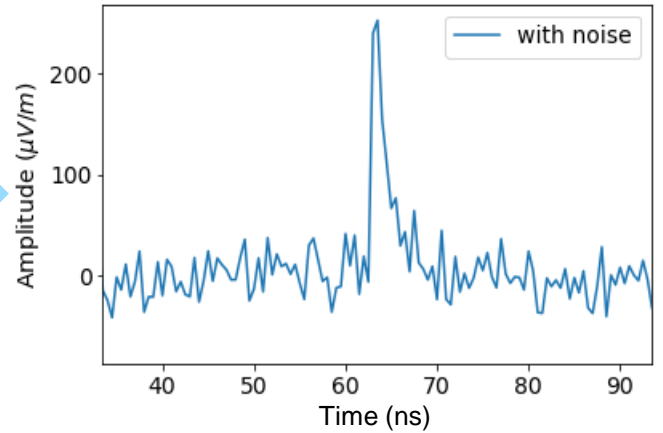
ZHAireS Simulations (Alvarez-Muñiz et al. 2011)

➔ Outputs: Traces $E_x(t)$, $E_y(t)$, $E_z(t)$ at each antenna

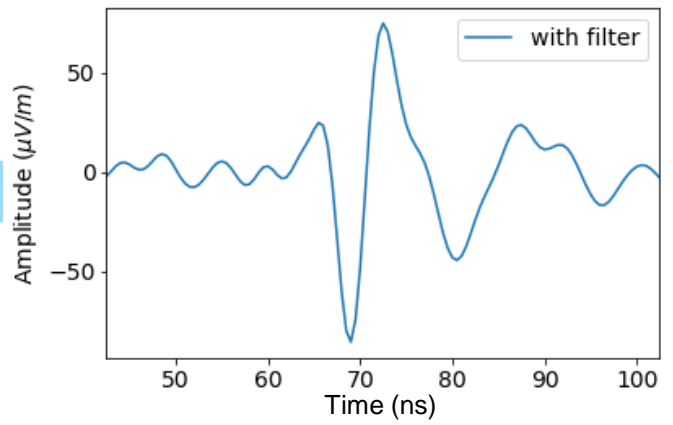
Account for experimental detection effects



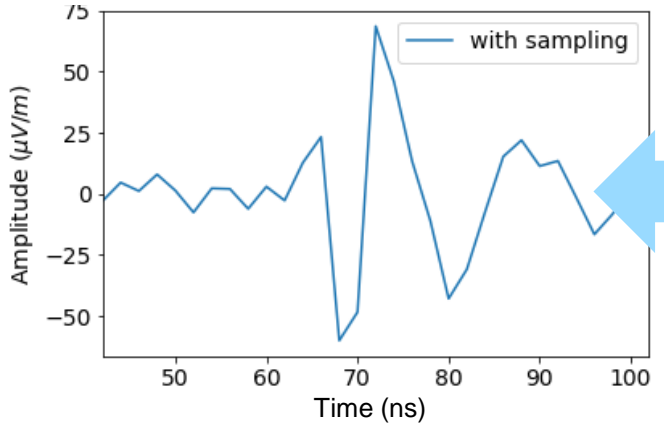
➔ gaussian noise (rms: 20 µV/m)



➔ filtering (50-200 MHz)



➔ sampling (2ns)



Shower plane

- Outputs of the simulations:
 $E_x(t)$, $E_y(t)$, $E_z(t)$
- We want to derive $E_v(t)$,
 $E_{v \times B}(t)$, $E_{v \times v \times B}(t)$

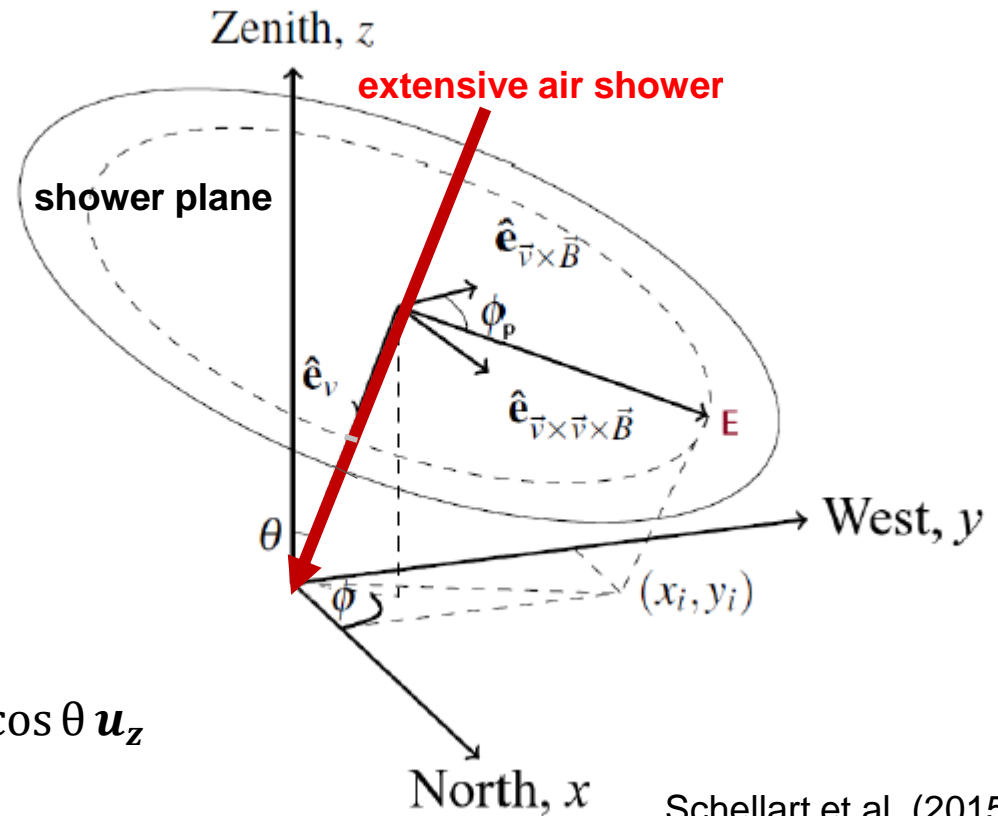
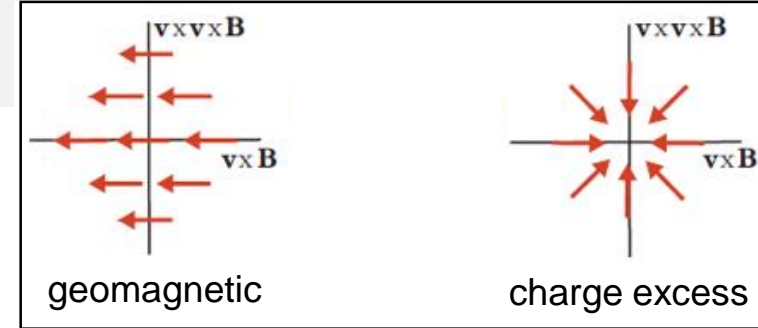
i : inclination of the magnetic field

θ : zenith angle

ϕ : azimuth of the shower

$$\mathbf{u}_B = \cos i \mathbf{u}_x - \sin i \mathbf{u}_z$$

$$\mathbf{u}_v = \sin \theta \cos \phi \mathbf{u}_x + \sin \theta \sin \phi \mathbf{u}_y + \cos \theta \mathbf{u}_z$$



Schellart et al. (2015)

We can derive $\mathbf{u}_{v \times B}$ and $\mathbf{u}_{v \times v \times B}$ from \mathbf{u}_v and \mathbf{u}_B and thus $E_v(t)$, $E_{v \times B}(t)$ and $E_{v \times v \times B}(t)$

Stokes parameters

- Stokes parameters I, Q, U, V: standard method to reconstruct the polarisation (Schoorlemmer 2012)
- $x_i = E_{v \times B}(t_i)$, $y_i = E_{v \times v \times B}(t_i)$,
- \hat{x}_i, \hat{y}_i , Hilbert transform of x_i, y_i , i.e., extension of the traces in the complex domain

Stokes parameters

100% Q	100% U	100% V
<p>+Q</p> <p>Q > 0; U = 0; V = 0 (a)</p>	<p>+U</p> <p>Q = 0; U > 0; V = 0 (c)</p>	<p>+V</p> <p>Q = 0; U = 0; V > 0 (e)</p>
<p>-Q</p> <p>Q < 0; U = 0; V = 0 (b)</p>	<p>-U</p> <p>Q = 0; U < 0; V = 0 (d)</p>	<p>-V</p> <p>Q = 0; U = 0; V < 0 (f)</p>

$$I = \frac{1}{n} \sum_{i=1}^n (x_i^2 + \hat{x}_i^2 + y_i^2 + \hat{y}_i^2) = |E_{v \times B}|^2 + |E_{v \times v \times B}|^2$$

$$Q = \frac{1}{n} \sum_{i=1}^n (x_i^2 + \hat{x}_i^2 - y_i^2 - \hat{y}_i^2) = |E_{v \times B}|^2 - |E_{v \times v \times B}|^2$$

$$U = \frac{2}{n} \sum_{i=1}^n (x_i y_i + \hat{x}_i \hat{y}_i)$$

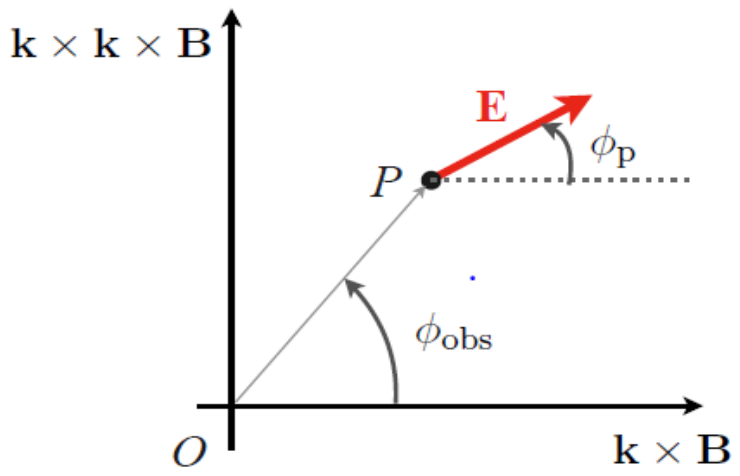
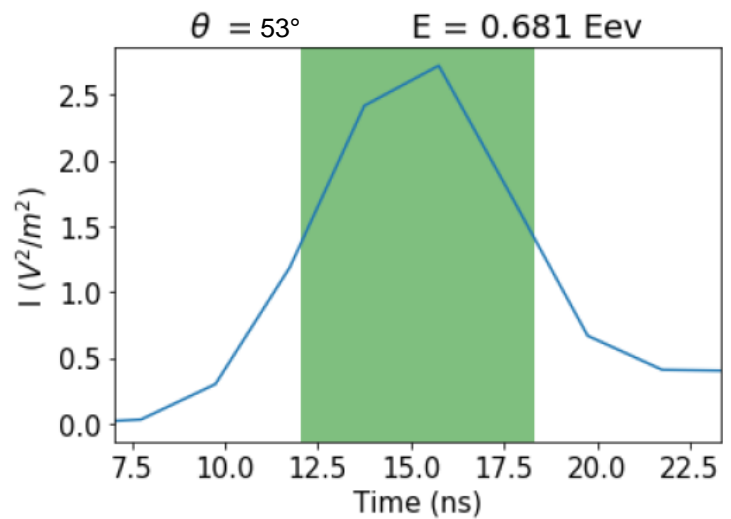
$$V = \frac{2}{n} \sum_{i=1}^n (\hat{x}_i y_i - x_i \hat{y}_i)$$

Reconstruction of the polarisation

- We have to define a time window over which we average the traces
- Stokes parameter I: Related to the total intensity of the traces

Time window: Fwhm of the I parameter

ϕ_p : angle between the polarisation and the $v \times B$ direction



$$\phi_p = 0.5 \tan^{-1} \frac{U}{Q}$$



$$E_{v \times B} = \sqrt{\langle I \rangle} \cos \phi_p$$

$$E_{v \times v \times B} = \sqrt{\langle I \rangle} \sin \phi_p$$

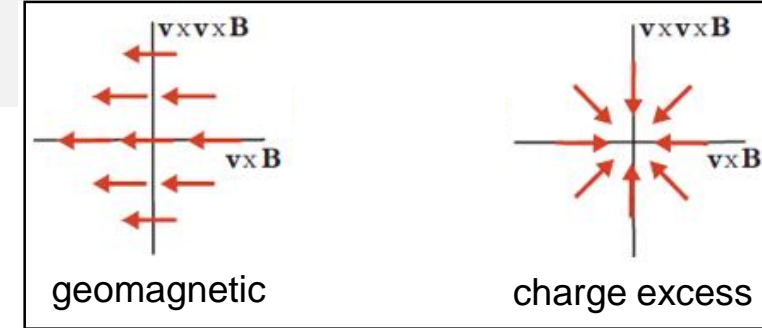
Reconstruction of the polarisation

Total polarisation

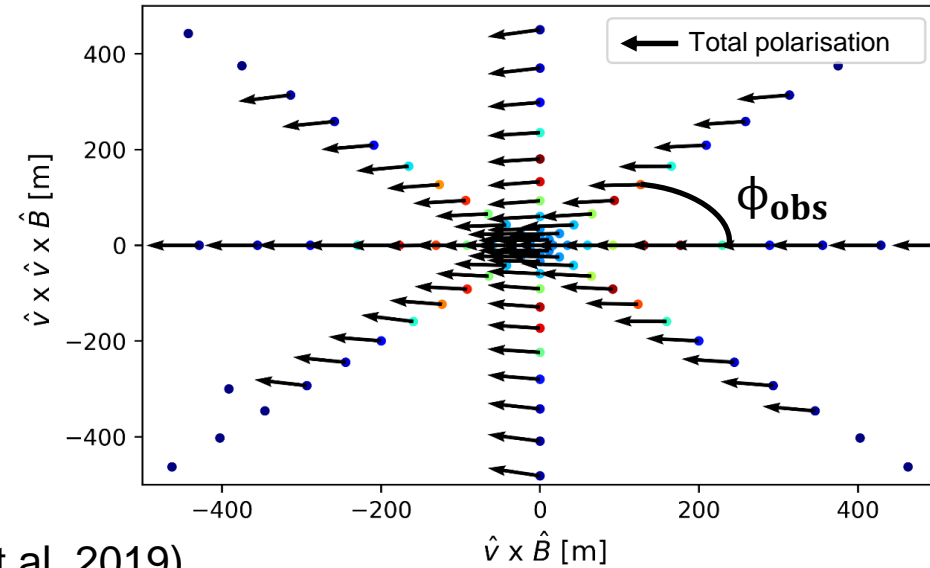
Various methods to reconstruct the polarisation:
absolute value, max value, Stokes parameters...

Total polarisation essentially along $-\mathbf{v} \times \mathbf{B}$

Dominant geomagnetic emission



Zenith = 53.31°, Energy = 3.981 Eev



Separation of each mechanism: (Huege et al. 2019)

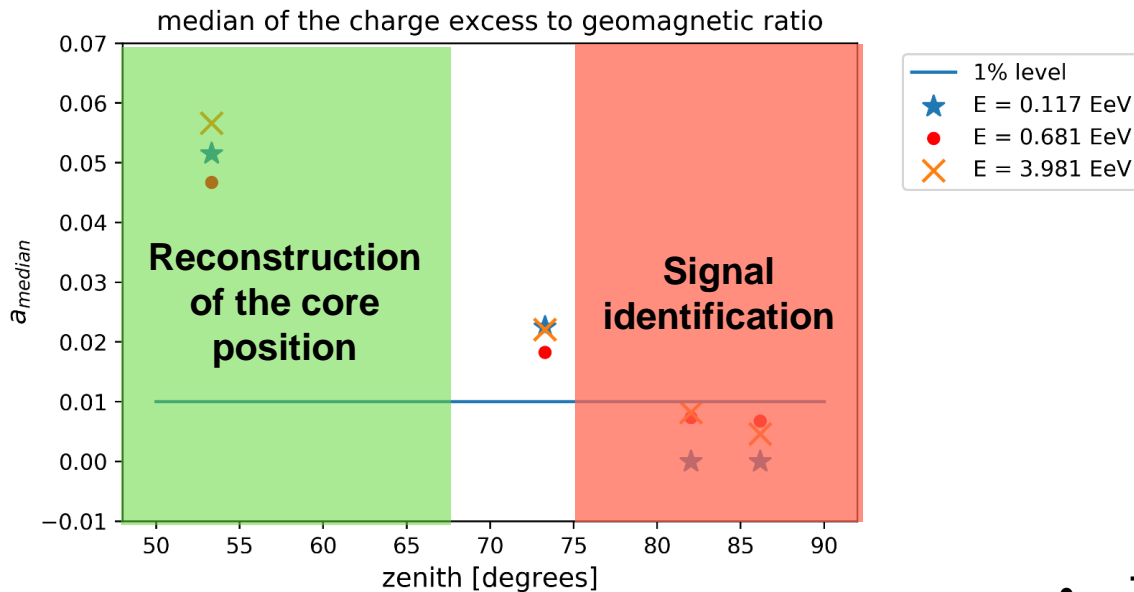
$$E_{ce} = \frac{E_{\mathbf{v} \times \mathbf{v} \times \mathbf{B}}}{|\sin \phi_{\text{obs}}|} \quad E_{\text{geo}} = E_{\mathbf{v} \times \mathbf{B}} - E_{\mathbf{v} \times \mathbf{v} \times \mathbf{B}} \frac{\cos \phi_{\text{obs}}}{|\sin \phi_{\text{obs}}|}$$

- Signatures to identify the radio signal
- Reconstruction of the air shower core position

Signal identification

Ratio of the amplitude of each mechanism

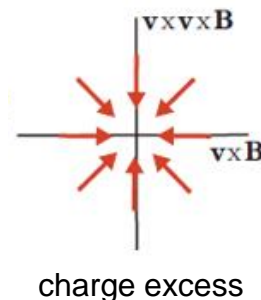
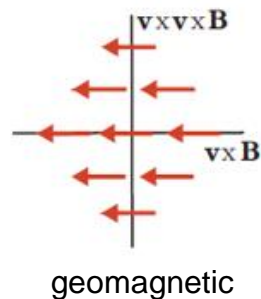
For different simulations



$$a = \sin \alpha \frac{E_{\text{charge excess}}}{E_{\text{geomagnetic}}}$$

Ratio below 1% for inclined air showers

Dominant contribution of the geomagnetic emission for inclined showers



- Total field orthogonal to B
- Strong signature of the radio signal visible directly at the antenna level
- Could be implemented in the trigger hardware of GRAND antennas

Shower core reconstruction

$$a = \sin \alpha \frac{E_{\text{charge excess}}}{E_{\text{geomagnetic}}}$$

The ratio drops to 0 at the core

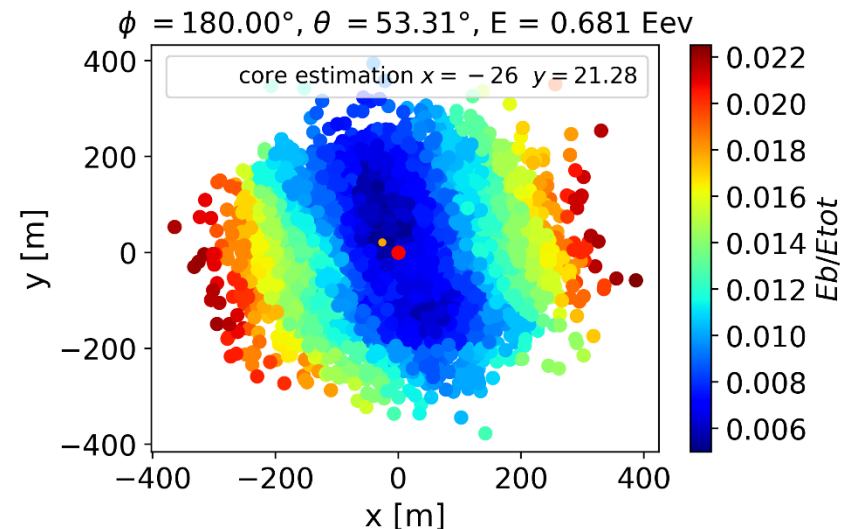
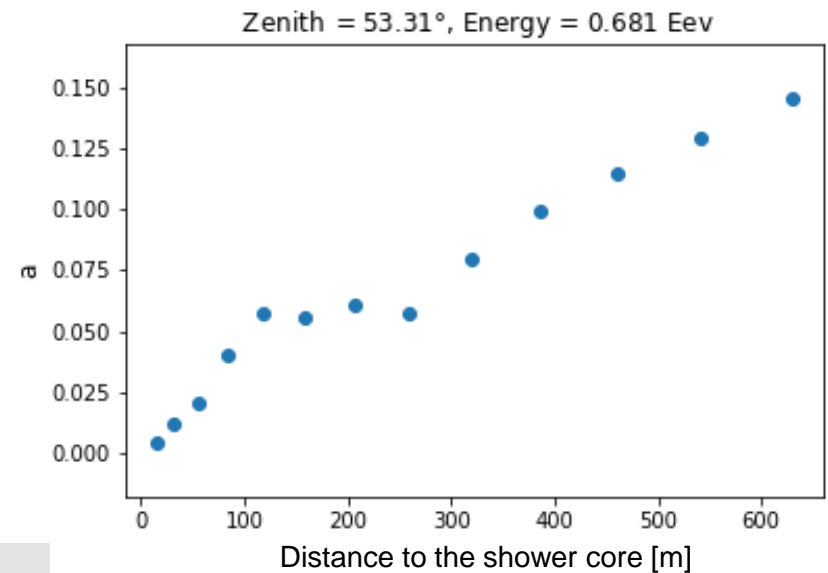
Increase with the distance to the core

Estimation of the shower core as the position that minimizes the ratio

Method

$$E_b = \mathbf{E}_{tot} \cdot \mathbf{u}_B$$

- In fact, we used the fraction of the total electric field along \mathbf{B} (E_b/E_{tot})
- For several positions we compute the mean ratio measured by the 20 closest antennas
- Core estimation: position with the lowest measurement



Still a preliminary work, but promising for showers with $\theta \leq 55^\circ$

Conclusion

Aim: Understanding the origin of ultra-high energy cosmic rays

- Multi-messengers approach to tackle this challenge
- Detection of the radio signal from extensive air showers induced by UHE astroparticles

GRANDProto300: Prototype of 300 antennas for the detection of UHE astroparticles

- Identification of the radio signal among the noise
- Reconstruction of the shower parameters

Results:

- Electric field orthogonal to \mathbf{B} for inclined showers
- The charge excess to geomagnetic ratio increases with distance to the core

