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The structure of hydrodynamic γ -ray burst jets

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ABSTRACT

After being launched, GRB jets propagate through dense media prior to their breakout. The jet-medium interaction results in the formation of a complex structured outflow, often referred to as a "structured jet". The underlying physics of the jet-medium interaction that sets the post-breakout jet morphology has never been explored systematically. Here we use a suite of 3D simulations to follow the evolution of *hydrodynamic* long and short gamma-ray bursts (GRBs) jets after breakout to study the post-breakout structure induced by the interaction. Our simulations feature Rayleigh-Taylor fingers that grow from the cocoon into the jet, mix cocoon with jet material and destabilize the jet. The mixing gives rise to a previously unidentified region sheathing the jet from the cocoon, which we denote the *jet-cocoon interface* (JCI). IGRBs undergo strong mixing, resulting in most of the jet energy to drift into the JCI, while in sGRBs weaker mixing is possible, leading to a comparable amount of energy in the two components. Remarkably, the jet structure (jet-core plus JCI) can be characterized by simple universal angular power-law angular distribution, and disfavors Gaussian jets. At larger angles, where the cocoon dominates, the structure is more complex. The mixing shapes the prompt emission lightcurve and implies that typical IGRB afterglows are different from those of sGRBs. Our predictions can be used to infer jet characteristics from prompt and afterglow observations.

Key words: gamma-ray burst — hydrodynamics — instabilities — methods: numerical – relativity



I-Introduction

- Prompt energetics \rightarrow collimated jet
- 2 populations of GRBs : IGRBs & sGRBs



Unified structure of GRB formation (Teboul & Piran+17)



I-Introduction

2 sites control jet structure :

- Launching site (initial acceleration and collimation)
- Interaction with surrounding medium during propagation

3D hydrodynamical simulations :

Turbulent mixing \rightarrow symmetry breaking \rightarrow Lateral energy dissipation

Results:

- Direct correlation mixing/lateral structure
- Post-breakout structure with 3 components (jet, JCI, cocoon)

II- Models

lGRB Model	$L_j[10^{50} {\rm erg \ s}^{-1}]$	$\theta_{j,0} = 0.7 \Gamma_0^{-1}$	<i>u</i> ∞, max	$M_{\star}[{ m M}_{\odot}]$	$ ho_{\star}(r)$	$t_b[s]$	$t_e[s]$
$Lc^{\dagger 1}$ (canonical)	1.0	0.14	500	10	$ ho_0 (r/r_0)^{-2} x^3$	20	68
Lw (wide)	1.0	0.18	400	10	$ ho_0(r/r_0)^{-2}x^3$	23	41
<i>Ln</i> (narrow)	1.0	0.07	1000	10	$ ho_0(r/r_0)^{-2}x^3$	13	24
<i>Lp</i> (powerful)	5.0	0.14	500	10	$ ho_0(r/r_0)^{-2}x^3$	8	36
<i>Lsd</i> (steep ρ profile)	1.0	0.14	500	2.5	$ ho_0(r/r_0)^{-2.5}x^3$	8	16
Lnp (narrow powerful)	7.0	0.07	1000	10	$ ho_0(r/r_0)^{-2}x^3$	6	33
$Lvp^{\dagger 2}$ (very powerful)	16	0.14	540	10	$ ho_0(r/r_0)^{-2}x^3$	5	16
<i>Llh</i> (low h)	1.0	0.14	100	10	$ ho_0(r/r_0)^{-2}x^3$	13	43
Lvwlh (very wide low h)	1.0	0.24	300	10	$ ho_0(r/r_0)^{-2}x^3$	27	55
Lvw (very wide)	1.0	0.24	500	10	$\rho_0(r/r_0)^{-2}x^3$	28	69
sGRB Model	$L_j[10^{50} { m erg s}^{-1}]$	$\theta_{j,0} \!= 0.7 \Gamma_0^{-1}$	$u_{\infty,\max}$	$M_{ce}[{ m M}_{\odot}]$	$\rho_*(r,\theta) [\mathrm{g \ cm^{-3}}]$	$t_d; t_b[s]$	$t_e[s]$
$S_{1}^{\dagger 3}$	1.4	0.07	200	0.04	$10^{22} (r/\text{cm})^{-2} \left(\frac{1}{4} + \sin^8\theta\right)$	0.2;0.4	1.0
$S_{2}^{\dagger 4}$	6.7	0.18	100	0.05	$5.5 \times 10^{34} (r/cm)^{-3.5}$	0.7;1.4	4.9
$S_{3}^{\tilde{7}5}$	0.3	0.14	500	0.05	$2.2 \times 10^{21} (r/cm)^{-2}$	0.6;1.4	3.6
S_4	10^{-3}	0.14	500	0.05	$2.2 \times 10^{21} (r/\text{ cm})^{-2}$	0.6;5.6	11.1

$$u_{\infty,max} = (\Gamma\beta)_{\infty,max} = \sqrt{h_0^2 \Gamma_0^2 - 1}$$

- t_d : Delay time
- t_b : Breakout time
- $t_{\rm e}$: Ejection time



- 4 regions :
- Jet
- Jet head
- Inner cocoon
- Outer cocoon

Rayleigh-Taylor instability :



Schwarzkopf+15



2D cut of the jet



Nuclear mushroom

Instability growth if

 $\widetilde{L_c} = \frac{\rho_j h_j \Gamma_j^2}{\rho_c h_c \Gamma_c^2} > 1$

Rayleigh-Taylor instability :



Schwarzkopf+15



Nuclear mushroom



2D cut of the jet

Richmeyer-Meshkov instability accelerates fingers growth

Relativistic Kevin-Helmoltz instability too slow





Intense mixing \rightarrow flatter energy distribution

More mixing ?

- Larger jet opening angle
- Higher jet specific enthalpy

→ sGRBs have stabler jets

- Lower jet luminosity
- Higher medium density

IV- The post-breakout structure



 $T = t - t_b$

 $E_i + E_{ICI} \simeq L_i T$

 $\lambda = E_i/L_iT$: mixing parameter In the JCI, $E_{iso} \propto \theta^{-\delta}$, $\delta = f(\lambda)$ 3 regions :

-Core

- ultra-relativistic:
$$u_{\infty} \ge \frac{1}{5} u_{\infty,max}$$

- narrow core angle: $\theta_j = \left(\frac{1}{3} - \frac{1}{5}\right) \theta_{j,0}$

Cocoon -

-

- newtonian:
$$u_{\infty} \leq 3$$

- large angles: $\theta \geq \theta_c \simeq 0.3 \ rad$
- Energy: $E_c \simeq L_i t_b$ -
- Jet-Cocoon Interface (JCI)
 - mildly relativistic -
 - $\theta_i \leq \theta \leq \theta_c$ intermediate angles: -



Energy distribution of the first slice of matter that breaks out, at various times

Energy distribution of matter between R_{\star} and 2 R_{\star} at various times

Overall energy distribution when the jet head reaches 10 R_{\star} , for different models

Is the post-breakout structure strongly asymmetric?



Energy distribution, sliced at different angles ϕ





Angular energy distribution at various times



Angular energy distribution, different models





Angular energy distribution at various times





Fit at a given time, with 3 distinct components



Fit at a given time, with 3 distinct components

$$E_{iso}(\lambda) \simeq \begin{cases} \lambda E_0 &, \text{ if } \theta < \theta_j \\ \lambda E_0 \left(\frac{\theta}{\theta_j}\right)^{-\delta} &, \text{ if } \theta_j < \theta < \theta_c \\ \lambda E_0 \left(\frac{\theta_c}{\theta_j}\right)^{-\delta} e^{-f_c(\theta - \theta_c)} &, \text{ if } \theta_c < \theta \end{cases}$$



$$u_{\infty}\left(\frac{\theta}{\theta_{j}}\right) = \begin{cases} < u_{\infty,j} > & \text{, if } \theta < \theta_{j} \\ < u_{\infty,j} > \left(\frac{\theta}{\theta_{j}}\right)^{-p_{u}} & \text{, if } \theta > \theta_{j} \end{cases}$$

Pure angular structure until $~\theta \sim 20~^\circ$

$$E_{iso}(\lambda) \simeq \begin{cases} \lambda E_0 &, \text{ if } \theta < \theta_j \\\\ \lambda E_0 \left(\frac{\theta}{\theta_j}\right)^{-\delta} &, \text{ if } \theta_j < \theta < \theta_c \\\\ \lambda E_0 \left(\frac{\theta_c}{\theta_j}\right)^{-\delta} e^{-f_c(\theta - \theta_c)} &, \text{ if } \theta_c < \theta \end{cases}$$

$$u_{\infty}\left(\frac{\theta}{\theta_{j}}\right) = \begin{cases} < u_{\infty,j} > & \text{, if } \theta < \theta_{j} \\ < u_{\infty,j} > \left(\frac{\theta}{\theta_{j}}\right)^{-p_{u}} & \text{, if } \theta > \theta_{j} \end{cases}$$

At later times, mixing drops:

- more energy in the core jet (λ increases)
- steeper energy decrease slope (δ increases)
- Steeper velocity gradient (p_u increases)

IGRB Model	$\lambda_4~(\lambda_{10})$	$\delta_4 (\delta_{10})$	$f_{c,4}(f_{c,10})$	$p_{u,4}(p_{u,10})$
Lc	0.1 (0.13)	1.2 (1.8)	1.6 (1.5)	2.0 (2.4)
Lw	0.09	0.8	1.3	2.1
Ln	0.18	1.8	2.6	1.8
Lp	0.15 (0.19)	1.7 (2.2)	2.8 (2.1)	1.9 (2.3)
Lsd	0.16	1.7	2.5	2.1
Lnp	0.27 (0.14)	2.2 (1.9)	2.5 (2.5)	2.3 (2.5)
Lvp	0.19	1.2	1.7	2.0
Llh	0.13 (0.25)	1.4 (2.0)	1.6 (1.9)	1.9 (1.9)
Lvwlh	0.08 (0.11)	0.7 (1.3)	0.9 (1.4)	2.1 (2.6)
Lvw	0.06 (0.11)	0.7 (1.1)	1.0 (1.4)	2.1 (2.7)

Similar as IGRBs with less dense surrounding medium

- ightarrow more collimated & stabler jets
- \rightarrow higher λ , δ
- \rightarrow cocoon extending to smaller angles



Overall energy distribution when the jet head reaches 10 R *, for different models



(b) Angular energy distribution, different models (c) 3-segments fitting

lGRB Model	$\lambda_4 \; (\lambda_{10})$	$\delta_4 \; (\delta_{10})$	$f_{c,4}(f_{c,10})$	$p_{u,4}(p_{u,10})$
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sGRB Model	λ	δ	f_c	p_u
S_1	0.4	3.1	4.5	2.7
S_2	0.43	3.5	3.7	2.4
S_3	0.38	3.2	4.1	3.2

A model with lower luminosity :



 \rightarrow Closer to IGRBs shapes

IV- The low-luminosity sGRBs problem

Breakout time is 10x too long !

Possible explanations:

- Ejected mass may be too high
- Ejecta might be highly anisotropic

With this we can produce sGRBs but might not account for r-process elements in the Universe...

Another explanation:

- Weakly magnetized jet \rightarrow fewer instabilities \rightarrow faster propagation \rightarrow more luminosity

IV- Gaussian or power-law angular profile ?



"As expected these distributions are incompatible with gaussian fits"

IV-2D vs. 3D



2D hydrodynamic simulation

3D hydrodynamic simulation

V- Emission

Prompt :

- light curves, spectra with temporal evolution
- some jets may not live long enough to produce a GRB
- higher mixing :
 - \rightarrow lower radiative efficiency
 - ightarrow high temporal efficiency variability
 - ightarrow alteration of off-axis emission

Afterglow :

- long and short GRBs : different light curves
- if mixing is weak (like in most sGRBs) :
 - \rightarrow if $\theta_{obs} < \theta_j$, afterglow similar to top-hat
 - \rightarrow if $\theta_{obs} > \theta_i$, afterglow peaks at late time, with a more complex rising phase

(hump & peak, 2 peaks)

- If mixing is high (most likely in IGRBs) :
 - \rightarrow if $\theta_{obs} < \theta_i$, early profile similar to top-hat, but decrease shallower
 - \rightarrow if $\theta_{obs} > \theta_j$, rise, peak, shallow decrease, fast decrease

VI- Summary

- Rayleigh-Taylor instabilities between cocoon and jet pre-breakout → mixing Mixing → Transition layer : JCI
- In IGRBs, after T = t_b , uniform energy distribution for jet in the star In sGRBs, less mixing \rightarrow more energy in the core
- Prior to breakout, almost all injected energy → cocoon
 After breakout, cocoon energy constant, injection in core & JCI
- IGRBs: more energy in the JCI of due to higher λ , larger JCI
- GRB distribution of E_{iso} and velocity follow a power law after the core
- 2D simulations of limited accuracy
- The presence of a magnetic field can stabilize the jet